ARCTIC CURVES BEYOND THE ARCTIC CIRCLE

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Pécs - Aug 2017

A PICTURE FROM THE BOOKLET



Figure 4: Picture by Vasarely

THE ARCTIC CIRCLE Jockusch-Propp-Shor



THE ARCTIC CIRCLE Jockusch-Propp-Shor

THE ARCTIC CIRCLE Jockusch-Propp-Shor



RANDOM TILING OF A SQUARE



dimer configuration: perfect matching on a graph (planar, doubly periodic, bipartite) ults hold for other (bipartite periodic planar) graphs



Dimer model on \mathbb{Z}^2 $\mathbb{Z}^2 \longleftrightarrow \text{dominos}$ P(z,w) = 1 + z + w - zw

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Dimer model on \mathbb{Z}^2 $\mathbb{Z}^2 \xrightarrow{} \operatorname{dominos}$ $\operatorname{pomino} \operatorname{tiling} \operatorname{model}_w$

dimer configuration: perfect matching on a graph (planar, doubly periodic, bipartite) ilts hold for other (bipartite periodic planar) graphs





honeycomb lattice \longleftrightarrow Dimer model on \mathbb{Z}^2 \rightarrow dominos ino tiling_model

 $\mathbf{N}_{\mathbf{r}}$

rhombus tilings

THE ARCTIC CIRCLE AGAIN Cohn-Larsen-Propp



THE ARCTIC CIRCLE AGAIN Cohn-Larsen-Propp



LOZENGETILINGS



LIMIT SHAPE THEOREM Cohn-Kenyon-Propp

For a fixed boundary height, as mesh size →0, with probability tending to 1 a random surface lies close to a deterministic surface, called *limit shape*

The limit shape is described by a variational principle



$$\begin{split} h: \Omega \to \mathbb{R} & \text{Lipschitz} \\ \min_{h} \int_{\Omega} \sigma(\nabla h), & \nabla h \in N \\ h_{|\partial\Omega} = h_0 \end{split}$$

analytic, strictly convex surface tension in the interior

degenerates on the boundary

SURFACE TENSION Cohn-Kenyon-Propp

Temperly-Fisher

Kasteleyn dimer model "exactly solvable" determinental formulae

$$\sigma(\mathbf{s}, \mathbf{t}) = -\frac{1}{\pi} (L(\pi \mathbf{s}) + L(\pi \mathbf{t}) + L(\pi(\mathbf{I} - \mathbf{s} - \mathbf{t}))) = -\mathbf{I}/\pi \operatorname{Vol}_{\mathbb{H}^3} \Sigma_{\alpha, \beta, \gamma}$$
$$L(\theta) = -\int_0^\theta \log(2\sin x) dx \quad \text{Lobachevsky function}$$



Limit shapes and the complex Burgers equation

by

RICHARD KENYON

Brown University Providence, RI, U.S.A. Andrei Okounkov

Princeton University Princeton, NJ, U.S.A.

U liquid region: where the Euler-Lagrange eq. makes sense
 $(h \text{ is } C' \text{ and } \nabla h \in \text{int } N)$ $h \in C^{\infty}(U)$ non-linear $\operatorname{div}(\nabla \sigma \circ \nabla h) = 0$ Morrey



non-linear



sense $h \in C^{\infty}(U)$ Morrey

Don'i jump, thínk conformal ínvaríance of height fluctuations ! (Kenyon-Okounkov)

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U liquid region: where the Euler-Lagrange eq. makes sense (h is C' and $\nabla h \in int N$) $h \in C^{\infty}(U)$ non-linear Morrey $\operatorname{div}\left(\nabla\sigma\circ\nabla h\right)=\mathbf{0}$ quasilinear $(\mathbf{x},\mathbf{y}) \in \mathbf{U} \rightsquigarrow \nabla \mathbf{h}(\mathbf{x},\mathbf{y}) = (\mathbf{s},\mathbf{t}) \rightsquigarrow \mathbf{f}(\mathbf{x},\mathbf{y}) \in \mathbb{D}$ **Kenyon-Okounkov:** "intrinsic coordinates" **7** $f: U \to \mathbb{D}$ $\bar{\partial}f = f \partial f$ harmonic S

l-s-t



E-L IN INTRINSIC COORDINATES $(s,t) \rightsquigarrow z \in \mathbb{C}_+$ $z = M \ddot{o} bius(f)$ $\nabla h = (\mathbf{s}, \mathbf{t})$ $\sigma_{s} = \log \frac{\sin \pi s}{\sin \pi (1 - s - t)} = \log |z|$ $\sigma_t = \log \frac{\sin \pi t}{\sin \pi (1 - s - t)} = \log |1 + z|$ $\nabla \sigma \circ \nabla h = (\log |\mathbf{z}|, \log |\mathbf{I} + \mathbf{z}|)$ $(\mathbf{x},\mathbf{y}) \in \mathbf{U} \mapsto \mathbf{z}(\mathbf{x},\mathbf{y})$ $\operatorname{\mathsf{Re}}\left(\frac{z_x}{z}+\frac{z_y}{1+z}\right)=\operatorname{\mathsf{div}}\left(\nabla\sigma\circ\nabla h\right)=0$ $h_{x} = \frac{1}{\pi} \arg(-(1+z))$ $h_{y} = -\frac{1}{\pi} \arg z$ $h_{xy} = h_{yx}$ $h_{xy} = h_{yx}$ In terms of $f: \overline{\partial} f = f \partial f$

ALGEBRAIC CURVES

Kenyon-Okounkov

frozen boundary = free boundary







polygonal boundary contour in coordinate directions

algebraic frozen boundary + frozen facets

deformation argument through volume constraint

"It would be interesting to prove ..."

ON THE VARIATIONAL PROBLEM deSilva-Savin

- - $\min_{h} \int_{\Omega} \sigma(\nabla h)$

 $\begin{array}{ll} \text{continuity} & \Gamma(\nabla h) & \text{is continuous} \\ & \Gamma\colon N\to \mathbb{S}^2 & \text{one-point compactification of int } N \\ & \partial N\leftrightarrow \{p\} \end{array}$

THE BELTRAMI EQUATION $\bar{\partial}\mathbf{f} = \mathbf{f}\partial\mathbf{f}$ $f \in W^{1,2}_{loc}(U)$ $U \subset \mathbb{C}$ $f \colon U \to \mathbb{D}$ $f \in C^{\infty}(U)$ bounded simply connected (or finitely connected) (self-improves) properness assumption: $f(z) \to \partial \mathbb{D}$, as $z \to \partial U$ $\nabla h \rightarrow \partial N$ at the boundary of U (critical boundary values will force formation of frozen facets) $f \in C(U)$

algebraic boundary

tangent relations

INVERSE MAP

$$f = B \circ g^-$$

 $g: \mathbb{D} \to U$ homeomorphism $\partial g = -B(z)\overline{\partial g}$ $B: \mathbb{D} \to \mathbb{D}$ analytic + proper \longrightarrow finite Blaschke product $h = g + B\overline{g}$ is holomorphic in \mathbb{D} $(|\mathbf{I} - |\mathbf{B}|^2)\mathbf{g} = \mathbf{h} - \mathbf{B}\overline{\mathbf{h}}$ $h(\zeta) = B(\zeta)\overline{h(\zeta)} = \frac{h(\zeta)}{B(\zeta)}$ $h(z) - B(z)\overline{h(z)} \rightarrow 0, \quad z \rightarrow \zeta \in \partial \mathbb{D}$ extend by reflection $h(z) = B(z)h(1/\overline{z}) \iff g(z) = g(1/\overline{z})$ $h(z) - h(I/\overline{z}) \to 0$, $|z| \to I$ $\overline{\partial}h = 0$ weakly across $\partial \mathbb{D}$ $\rightarrow h$ is a rational map of \mathbb{C}

UNIVALENT POLYNOMIALS $B(z) = z^d, \quad d \ge 2$ $h(z) = z^d \overline{h(1/\overline{z})}, \quad h(0) = 0 \qquad h(z) = \alpha z + \ldots + \overline{\alpha} z^{d-1}$ *h* is a self-reflective polynomial of degree *d*-*l* $g(z) = \frac{h(z) - z^d \overline{h(z)}}{|I - |z|^{2d}} \quad \begin{array}{l} \text{extends continously to the boundary} \\ g(\zeta) = p(\zeta), \quad \zeta \in \partial \mathbb{D} \end{array}$ $p(z) = h(z) - \frac{1}{d}z h'(z)$ $p(\partial \mathbb{D})$ non self-crossing curve \rightarrow *p* is univalent in \mathbb{D} $p'(z) = (\mathbf{I} - \frac{\mathbf{I}}{d})\alpha + \ldots + (\mathbf{I} - \frac{\mathbf{I}}{d})\bar{\alpha}z^{d-2}$ \rightarrow all d-2 critical points are on $\partial \mathbb{D} \rightarrow z^{d-2}\overline{p'(1/\overline{z})} = p'(z)$



RATIONAL PARAMETRIZATION

 $g(z) = \frac{h(z) - B(z)\overline{h(z)}}{1 - |B(z)|^2} = \frac{\frac{h(z)}{B(z)} - \overline{h(z)}}{\frac{1}{B(z)} - \overline{B(z)}} = \frac{\frac{h(z)}{B(z)} - \frac{h(1/\overline{z})}{B(1/\overline{z})}}{\frac{1}{B(z)} - \frac{1}{B(1/\overline{z})}}$ $r = \frac{(h/B)'}{(1/B)'} \qquad \overline{r(1/\overline{z})} = \frac{h'}{R'}$ $\rightarrow r(\zeta), \quad z \rightarrow \zeta \in \partial \mathbb{D}$ zr'(z) self-reflective with respect to B(z) $\frac{B(z)}{z^2} \frac{1}{r'(1/\overline{z})} = r'(z) \qquad \text{tangent vector} \quad i\zeta r'(\zeta) = iA(\zeta)\sqrt{B(\zeta)}$ real, sign changes d/2 - 1/2 # cusps = 1at cusps boundary is algebraic & locally convex except at d-2 cusps $f = B \circ g^{-1}$ continuous up to the boundary 3d-3 real parameters

BOUNDARY RECONSTRUCTION



MULTIPLY CONNECTED CASE $f = B \circ g^{-1}$

 $\mathbb{D} \rightsquigarrow \mathcal{D} \quad \text{circle domain} \quad B \colon \mathcal{D} \to \mathbb{D} \quad \text{analytic and proper}$

B and h continuous up to the boundary as before

$$h(\zeta) = B(\zeta)\overline{h(\zeta)} = \overline{\frac{h(\zeta)}{B(\zeta)}}, \quad \zeta \in \partial \mathcal{D}$$

 $\begin{array}{l} (h,h/B) \\ (B,\overline{1/B}) \end{array} \text{ are meromorphic functions on the Schottky double } \hat{\mathcal{D}} \\ j \text{ antiholomorphic involution} \end{array}$

$$\tilde{r} = \left(\frac{h'}{B'}, \frac{\overline{(h/B)'}}{(1/B)'}\right) \qquad r = \overline{\tilde{r} \circ j}$$

 $g(\zeta) = r(\zeta), \quad \zeta \in \partial \mathcal{D}$

meromorphic functions $R(r, \tilde{r}) = 0$ $R(g(\zeta), \overline{g(\zeta)}) = 0$ algebraic boundary etc

Arctic curves of the octahedron equation

Philippe Di Francesco¹ and Rodrigo Soto-Garrido²



Arctic curves of the octahedron equation

Philippe Di Francesco¹ and Rodrigo Soto-Garrido²

 $P(u, v) = 603\,358\,073\,569\,688\,095\,393\,738\,000u^{14} + 1822\,971\,422\,522\,481\,873\,814\,304\,800vu^{13} + 302\,414\,835\,014\,281\,399\,576\,977\,600u^{13}$ $+7658\,013\,562\,515\,635\,323\,215\,886\,000v^2u^{12}+626\,386\,479\,045\,976\,264\,625\,165\,760vu^{12}+65\,648\,625\,922\,043\,130\,480\,407\,960u^{12}$ $+8502\,660\,801\,885\,990\,861\,442\,260\,800v^{3}u^{11}-4016\,291\,377\,989\,674\,598\,197\,523\,840v^{2}u^{11}-8955\,889\,812\,423\,159\,779\,663\,425\,824vu^{11}$ $+8631816024097405173283346160v^{2}u^{9}+6130342332781365103023636918vu^{9}+300038159586641951467587240u^{9}$ $+48\,101\,067\,368\,389\,417\,583\,947\,124\,400v^{6}u^{8}-8874\,912\,221\,343\,735\,420\,641\,284\,800v^{5}u^{8}-30\,642\,500\,612\,723\,566\,034\,952\,420\,120v^{4}u^{8}$ $+4089814979226593490453920400v^{3}u^{8}+2890833100949061663021542421v^{2}u^{8}-1537862122247709326673670200vu^{8}$ $-1276791684735224437145235252u^{8} + 165638160476224209249648000v^{7}u^{7} - 15185547478970793846022129920v^{6}u^{7}$ $+10\,619\,324\,480\,232\,243\,252\,222\,805\,440v^5u^7$ + 17 121 927 414 923 400 963 351 428 640 v^4u^7 - 4642 084 019 946 561 205 079 466 936 v^3u^7 $-4201\,308\,893\,745\,605\,384\,096\,673\,600v^2u^7$ + 97 128 658 780 698 750 571 038 384 vu^7 + 16 658 271 644 450 437 458 125 640 u^7 $+48\,101\,067\,368\,389\,417\,583\,947\,124\,400v^8u^6-15\,185\,547\,478\,970\,793\,846\,022\,129\,920v^7u^6-54\,696\,534\,109\,775\,129\,942\,931\,200\,352v^6u^6$ $+8498087480515562992290313440v^5u^6 + 20848735934263779279940738242v^4u^6 - 2928090072842649426783830400v^3u^6$ $- 1125\,942\,030\,946\,106\,640\,101\,862\,864 v^2 u^6 + 881\,693\,827\,811\,784\,667\,334\,364\,120 v u^6 + 410\,818\,358\,444\,129\,895\,320\,450\,118 u^6$ $+8664424609796383599417068000v^9u^5 - 8874912221343735420641284800v^8u^5 + 10619324480232243252222805440v^7u^5$ $+ 8498\,087\,480\,515\,562\,992\,290\,313\,440v^6u^5 - 16\,598\,910\,777\,434\,586\,615\,901\,305\,852v^5u^5 - 3118\,943\,690\,894\,703\,413\,913\,413\,040v^4u^5$ $+4436727620735139576883870032v^{3}u^{5}+378779090210933672213353800v^{2}u^{5}-894275420028329313474734772vu^{5}$ $-28\,143\,830\,188\,642\,461\,399\,955\,080u^{5}+24\,870\,815\,123\,962\,290\,558\,144\,794\,000v^{10}u^{4}-17\,028\,590\,764\,390\,279\,389\,912\,000v^{9}u^{4}$ $-30\,642\,500\,723\,566\,034\,952\,420\,120v^8u^4 + 17\,121\,927\,923\,400\,963\,351\,428\,640v^7u^4 + 20\,848\,735\,263\,779\,279\,940\,738\,242v^6u^4$ $-3118943690894703413913413040v^5u^4 - 6585025120215513060415620600v^4u^4 + 224576822600011254994156440v^3u^4$ $+730\,062\,356\,407\,169\,871\,489\,508\,026v^2u^4 - 87\,999\,348\,446\,432\,687\,845\,418\,760vu^4 - 39\,991\,576\,579\,826\,072\,416\,315\,884u^4$ $+4089814979226593490453920400v^8u^3 - 4642084019946561205079466936v^7u^3 - 2928090072842649426783830400v^6u^3$ $+4436727620735139576883870032v^5u^3 + 224576822600011254994156440v^4u^3 - 771752886154129578670446744v^3u^3$ $+54\,105\,975\,565\,681\,638\,845\,373\,840v^2u^3$ + 158 742 939 499 283 087 522 192 736 vu^3 + 3181 828 983 737 934 822 021 000 u^3 $+7658013562515635323215886000v^{12}u^2 - 4016291377989674598197523840v^{11}u^2 - 6515407606857043381218037200v^{10}u^2$ $+8631816024097405173283346160v^9u^2 + 2890833100949061663021542421v^8u^2 - 4201308893745605384096673600v^7u^2$ $-1125942030946106640101862864v^6u^2 + 378779090210933672213353800v^5u^2 + 730062356407169871489508026v^4u^2$ $+54\,105\,975\,565\,681\,638\,845\,373\,840v^{3}u^{2}-205\,856\,416\,682\,486\,477\,443\,753\,704v^{2}u^{2}-4541\,013\,871\,098\,771\,634\,821\,000vu^{2}$ $-417838190775940873949175u^{2} + 1822971422522481873814304800v^{13}u + 626386479045976264625165760v^{12}u$



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LIFE BEYOND THE ARCTIC CIRCLE

